

Rethinking the Macroscopic Presentation of the Second Law of Thermodynamics

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Abstract: The classical macroscopic presentation of the second law of thermodynamics is an elegant but abstract sequence of very specific thought experiments that utilize reversible processes occurring within heat engines operating between infinite temperature reservoirs. The length, specificity and complexity of this sequence may hamper the understanding of important concepts such as exergy and entropy. The pedagogical problems of this approach have been discussed, followed by an alternative presentation wherein second law concepts and formulations have been derived from thought experiments that use real, rather than imaginary processes. The thought experiments involve classifying heat transfer at any local point for any arbitrary process involving work-heat interactions into different categories, and then collecting terms for each category throughout the control volume in order to relate property changes to external heat transfer and/or work. They embrace the spatial non-uniformity present in any real process, are consistent with contemporary computational approaches, and can potentially serve as building blocks for the development of computational thinking in students. An assessment plan with limited sample size has been described. The primary purpose of this paper to interest other thermodynamics instructors in the proposed presentation so that the assessment can be performed with a large number of students

1. Pedagogical Problems with the Classical Presentation: The approximate sequence of the classical macroscopic presentation of second law concepts and results has not changed for more than a century. Figure 1 shows a schematic of the sequence of steps followed by engineering textbooks used in introductory thermodynamics courses over the last half-century¹⁻⁹, all based on the edifice constructed by Carnot, Clausius and Kelvin¹⁰. The sequence commencing with the Kelvin-Planck/Clausius statements of the Second Law and culminating with Exergy analysis is long, sometimes spanning more than 200 pages in recent textbooks¹, complex and completely based on imaginary reversible processes.

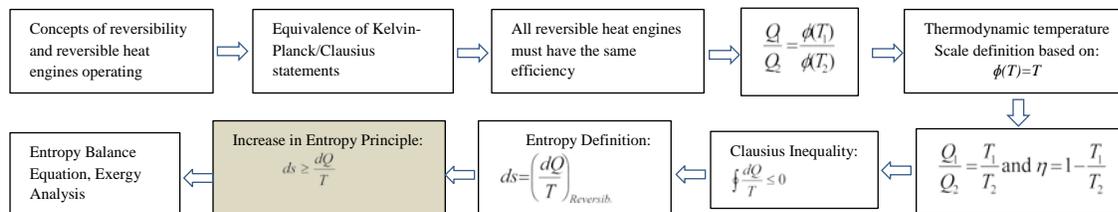


Figure 1. Schematic of the sequence of steps in the classical presentation of the second law in engineering textbooks over the last fifty years¹⁻⁹. The proposed presentation derives the principle of increase of entropy directly, for any general process involving work and/or heat transfer.

Although the classical presentation is stimulating in an abstract intellectual sense, it has a number of shortcomings from a pedagogical perspective:

a. Derivation is unrelated to application: Mathematical formulations for calculating exergy, entropy generation and irreversibilities follow from the Clausius inequality, all of which are derived from arguments that utilize imaginary reversible heat engines (RHEs) in imaginary situations. RHE efficiency in turn originates from a seemingly arbitrary choice of temperature function used to define the Thermodynamic Temperature Scale (TTS). Ultimately, all second law formulations are derived using infinitely slow reversible processes during which all properties are spatially uniform. The insight gained by following and understanding the derivation is not directly transferable to the second law analysis of any real system.

b. Specific-to-general approach: The derivations are undertaken with specific devices (heat engines) and processes (reversible processes) but students are expected to apply the second law to general problems that do not use these particular devices or processes, e.g. exergy analysis of a real (irreversible) fuel cell. This specific-to-general approach is an exception to the general pedagogical practice of deriving results for a general situation that is then applied to specific cases.

43 c. Entropy is an abstract concept: Determining the entropy change between two states requires traversing
44 an imaginary reversible path between them. Entropy might have been an abstract concept in the twentieth
45 century but it is defined, understood and used as a measure of dispersion in real systems, in many
46 contemporary fields such as data mining and information theory. The proposed derivation is consistent
47 with this modern approach; Entropy is defined at an infinitesimal point for a real process so that entropy
48 generation is understood fundamentally in terms of dispersion of heat resulting from spatial non-
49 uniformity.

50 d. Irreversibility is poorly understood: The classical presentation precedes the development of
51 computational approaches that describe spatial non-uniformities. All derivations require spatially uniform
52 (and therefore infinitely slow) processes. Irreversible processes are simply as processes that are *not*
53 reversible. If future engineers are going to design devices with high second law efficiency by minimizing
54 irreversibilities, they need to understand irreversibilities in terms of spatial non-uniformity of processes
55 and properties.

56 Outside the universe of engineering textbooks, the second law has been expressed and formulated in
57 many different ways for different audiences, e.g. works by Morales¹¹, Macdonald¹², Muschik¹³,
58 Thomsen¹⁴ and Baerlein¹⁵. A discussion of different second-law approaches can be found in a review
59 paper by Muschik¹⁶. None of these approaches address the four pedagogical shortcomings listed above;
60 they are still based on RHE's operating between temperature reservoirs. Many Introductory physics
61 textbooks at the college level have modified their presentation of the second law by introducing entropy
62 from a molecular perspective, while using an abridged version of the sequence shown in figure 1 to
63 discuss only RHE's (exergy is generally not covered). Some introductory physics textbooks¹⁷⁻²⁰ skip the
64 Clausius theorem altogether, and derive RHE efficiency starting from $\Delta S=0$. Others derive the Clausius
65 theorem from RHE efficiency²¹, which is presented as the upper limit of efficiency (without the RHE
66 corollaries presented in almost all engineering textbooks) after being derived for an ideal gas.

67
68 The motivation for the current work is to address the four shortcomings listed above by deriving all
69 macroscopic second law results for any arbitrary real process involving heat and/or work transfer.
70

71 **2. Proposed Derivation of Second Law Formulations for any Arbitrary (Real) Process**

72 The derivation is divided into two parts; the Local Heat Category (LHC) equation introduced in this work
73 is presented first. It is then used to derive the standard second-law results found in introductory textbooks.

74 2.1 Local Heat Category (LHC) Equation

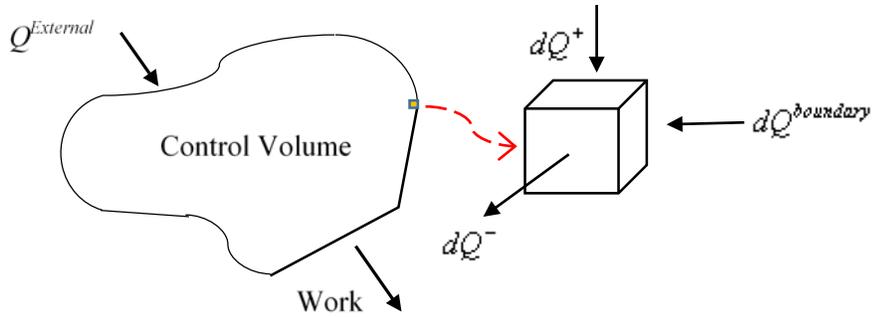
75 Consider heat transfer across the surfaces of an infinitesimally small volume inside a finite control
76 volume (CV) as shown in figure 1.

77 Energy transferred as heat at this infinitesimal point can be classified into three exclusive categories:

78 **a. Internal heat transfer from another internal point excluding the external source.** This positive heat
79 transfer term will be denoted by dQ^+ .

80 **b. Internal heat transfer to another internal point excluding the external sink.** This negative heat transfer
81 term will be denoted by dQ^- .

82 **c. External heat transfer from/to the external source/sink.** This can happen at points located along the
83 boundary of the CV, or through radiation to /from internal points. This term will be denoted by
84 $dQ^{boundary}$.



85
 86 **Fig 1.** The terms of the LHC equation shown at an infinitesimal point inside a finite CV. Heat transferred from and to other
 87 interior points are denoted by dQ^- and dQ^+ respectively. Heat transfer to or from external sources/sinks is denoted by $dQ^{boundary}$.
 88 For simplicity, each term has been shown to act across one face only. In general however, each term is comprised of flux from all
 89 faces, as per equation 2.

90

91 The net heat energy gained/lost at any point inside the CV is then given by:

92

$$93 \quad dQ = dQ^- + dQ^+ + dQ^{boundary} \quad (1)$$

94

95 This is the Local Heat Category (LHC) equation. Each term represents the infinitesimal amount of energy
 96 transferred across the surfaces of an infinitesimal volume over an infinitesimal time duration. For
 97 example:

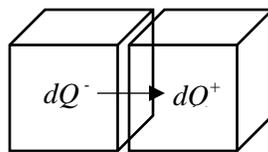
98

$$99 \quad dQ^- = \left(\sum_{i=1}^6 q_i''^- dA_i \right) dt \quad (2)$$

100

101 where $q_i''^-$ is the instantaneous negative heat flux across an infinitesimal surface of area dA_i . Heat
 102 energy rather than heat flux terms will be henceforth used because of simplicity, so equation (2) is not
 103 part of the derivation.

104



105
 106 **Fig 2.** The dQ^- term for the infinitesimal point on the left becomes the dQ^+ term for the point on the right. For simplicity, each
 107 term has been shown to act across one face only. In general the dQ^- term would be comprised of fluxes from multiple faces, as
 108 per equation (2), and disperse as part of the dQ^+ term for multiple points. When collected, the sum of both terms will still be zero,
 109 as shown by equation (3).
 110

111 Since dQ^- is internal by definition, every dQ^- term across the surface of every infinitesimal point must
 112 be part of a dQ^+ term elsewhere (also internal by definition), and vice-versa, as shown simplistically by
 113 figure 2. The sum of all dQ^- terms across all faces of all infinitesimal points must be equal in magnitude
 114 to the corresponding sum of dQ^+ terms. This is shown by equation (3):

115

116
$$\int_{CV} dQ^- + \int_{CV} dQ^+ = 0 \quad (3)$$

117
 118 Equation (3) simply says that the sum of positive internal heat transfer throughout the CV is equal to the
 119 sum of negative internal heat transfer. It can be integrated over a finite time duration:
 120

121
$$\int \int_{CV} dQ^- + \int \int_{CV} dQ^+ = 0 \quad (4)$$

122 Note that the double integral produces finite terms because the dQ terms are products of two differential
 123 quantities as per equation (2). Second law results will be derived for finite processes, so double integral
 124 will be used henceforth.

125 An important result follows directly from equation (4), and the Second Law statement that heat transfer
 126 can only occur from higher to lower temperature. Since the temperature at the internal source(s) of the

127 dQ^- term must exceed the temperatures at the locations corresponding to the dQ^+ term, $\left| \frac{dQ^-}{T} \right|$ must be

128 smaller than $\left| \frac{dQ^+}{T} \right|$. Therefore:

129
$$\int \int_{CV} \frac{dQ^-}{T} + \int \int_{CV} \frac{dQ^+}{T} \geq 0 \quad (5)$$

130
 131 because, the first term is positive while the second is negative. This result will be used later.

132 2.2 Derivation of $ds \geq \frac{dQ}{T}$ for any Arbitrary (Real) Process

133 Consider any arbitrary process involving external heat transfer to or from any CV as shown in figure 1.
 134 Multiple heat sources and/or sinks might exist and external work may/may not be done on/by the CV. If
 135 $q''^{External}$ is the instantaneous heat flux at any point on the surface of the CV, then the net external heat
 136 transfer is given by:
 137

138
$$Q^{External} = \int \int_A q''^{External} dA dt = \int \int_A dQ^{External} \quad (6)$$

139
 140 The external heat flux is integrated over the surface area of the CV, denoted by A . All of the external heat
 141 transfer must occur across the boundary of the CV. Therefore:
 142

143
$$\int \int_{CV} dQ^{boundary} = \int \int_A dQ^{External} \quad (7)$$

144
 145 where $dQ^{boundary}$ is the last term of the LHC equation, see equation (1). Note that integration of $dQ^{boundary}$
 146 over the CV implies integration over the surfaces of infinitesimal points as per equation (2). This leads to
 147 the second important result following equation (5):
 148

149
$$\int \int_{CV} \frac{dQ^{boundary}}{T} \geq \int \int_A \frac{dQ^{External}}{T} \quad (8)$$

150

151 The argument identical to the one used to obtain equation (5); the temperature of any boundary point
 152 inside the CV has to be lower than the external source temperature, or higher than the external sink
 153 temperature. For the latter case, the left-hand-side will be a smaller negative number than the right-hand-
 154 side.

155 Adding the two important results, equations (5) and (8):

156

$$157 \quad \int \int_{i \text{ CV}} \frac{dQ^-}{T} + \int \int_{i \text{ CV}} \frac{dQ^+}{T} + \int \int_{i \text{ CV}} \frac{dQ^{\text{boundary}}}{T} \geq \int \int_{i \text{ A}} \frac{dQ^{\text{External}}}{T} \quad (9)$$

158

159 Re-organizing terms and using the LHC equation (1), we obtain:

160

$$161 \quad \int \int_{i \text{ CV}} \frac{dQ^- + dQ^+ + dQ^{\text{boundary}}}{T} = \int \int_{i \text{ CV}} \frac{dQ}{T} \geq \int \int_{i \text{ A}} \frac{dQ^{\text{External}}}{T} \quad (10)$$

162

163 If the term $\frac{dQ}{T}$ is denoted by the variable dS :

$$164 \quad \int \int_{i \nabla} dS \geq \int \int_{i \text{ A}} \frac{dQ^{\text{External}}}{T} \quad (11)$$

165

166 It can be easily shown that the variable dS is a point function; therefore 'S' is a property that will be called
 167 entropy. Equation (11) is the familiar mathematical statement of the second law, derived from heat engine
 168 arguments in the classical presentation. The proposed approach reverses this specific-to-general approach,
 169 and uses equation (11) to derive all the mathematical results of the second law, including reversible heat
 170 engine (RHE) efficiency, as illustrated in the next sub-section 3.3.

171

172 2.3 Mathematical Results Following from the Equation (11)

173 For any cyclic process, the property change $\Delta S=0$ and equation (11) reduces to the Clausius Inequality:

174

$$175 \quad \oint_A \frac{dQ^{\text{External}}}{T} \leq 0 \quad (12)$$

176

177 It is evident that in order to convert heat into work, at least one heat sink would be required in order for
 178 the left-hand-side to be negative. This is consistent with the Kelvin-Planck statements of the second law.
 179 For the limiting case where the CV encloses a cyclic and reversible heat engine (RHE) operating between
 180 a single source and a single sink of constant temperature, equation (12) reduces to:

181

$$182 \quad \frac{\Delta Q^{\text{Source}}}{T^{\text{Source}}} + \frac{\Delta Q^{\text{Sink}}}{T^{\text{Sink}}} = 0 \quad (13)$$

183

184 This results in the familiar expression for the thermal efficiency of a RHE operating between two
 185 temperature reservoirs:

$$186 \quad \eta_{\text{reversible}} = \frac{\Delta Q^{\text{Source}} + \Delta Q^{\text{Sink}}}{\Delta Q^{\text{Source}}} = 1 - \frac{T^{\text{Source}}}{T^{\text{Sink}}} \quad (14)$$

187 The second law equation (11) can also be used to determine the exergy of a substance. If the CV encloses
 188 the substance (without a heat source), it can be seen that a heat sink will be required to achieve a change
 189 of state, i.e. non-zero ΔS . Equation (11) then integrates to:
 190

$$191 \quad \Delta S \geq \frac{\Delta Q^{Sink}}{T_o} \quad (15)$$

192 where ΔS corresponds to the change between current and dead state. Maximum work production will
 193 correspond to minimum heat rejection, i.e. the limiting equality corresponding to an imaginary reversible
 194 process:
 195

$$196 \quad \Delta Q_{reversible}^{Sink} = T_o \Delta S \quad (16)$$

197 The mathematical expressions for exergy of any closed or open system readily follow from equation (16)
 198 when combined with the first law. The ΔS term must include the entropy change of the flow terms if mass
 199 crosses the CV boundaries.
 200

201 **3. Pedagogical Implications of Proposed Derivation:** The proposed derivation makes it easier to
 202 understand (ir)reversibility, entropy, entropy generation and exergy destruction in real and arbitrary
 203 systems that are not heat engines. The only second law statement used for the derivation is that heat is
 204 transferred from higher to lower temperatures. Students understand this intuitively and can appreciate that
 205 everything else follows from this. Reversibility can be mathematically defined using the LHC equation as
 206 any process in which:
 207

$$208 \quad dQ^- = dQ^+ = 0 \quad (17)$$

209 at every point and every instant throughout the process. This is because every kind of irreversibility
 210 within a CV, including frictional dissipation, will ultimately result in irreversible internal heat transfer.
 211 The definition unites the different kinds of irreversibilities that are described in current textbooks as any
 212 process violating equation (17). In that case it is easy to see that every violation results in the loss of work
 213 potential. The differential amount of exergy lost when dQ^- becomes dQ^+ across temperature difference dT
 214 can be easily derived from equation (14):
 215

$$216 \quad dX = \frac{dQ^+}{T^2} dT \quad (18)$$

217 Students can then understand why reversible processes conserve work potential while irreversible
 218 processes do not. They can appreciate that dQ^+ and dQ^- terms will be non-zero for a real process, but
 219 minimizing them can reduce exergy destruction. They can intuitively understand that dQ^+ and dQ^- terms
 220 can be minimized by minimizing temperature gradients within the system, by designing processes that are
 221 spatially uniform. For example, the exergy destruction in any combusting system is greatly reduced if
 222 combustion occurs uniformly throughout the combustion chamber²².
 223

224 Entropy is defined at an infinitesimal point, see equation (11), and calculating entropy change does not
 225 require a parallel reversible process. Entropy generation can be understood to result from $\frac{dQ^+}{T}$ and $\frac{dQ^-}{T}$
 226 terms, and ultimately from spatial non-uniformity. A more detailed discussion of defining entropy in this
 227 manner can be found elsewhere²³, and is best suited to a graduate thermodynamics course.
 228

232
233 The derivation can also be used to reduce confusion between the different temperature scales. Unlike the
234 classical presentation which requires defining a Thermodynamic Temperature Scale (TTS), see figure 1,
235 the proposed derivation can be based on the more easily understood Ideal Gas Temperature Scale (IGTS).
236 The TTS (or the Kelvin scale) can then be derived from equation (14) to show that the IGTS coincides
237 with the TTS. Again, the reader is pointed elsewhere²³ for a more detailed discussion on this topic.
238

239 **4. Assessment:** A simple assessment method would be to compare two cohorts of students who have been
240 taught the classical and proposed presentations respectively, using a well-validated measurement tool.
241 One such tool is the Concept Inventory for Engineering Thermodynamics (CIET) developed by Vigeant
242 et al²³. Reliability data for the CIET was collected from 15 institutions nationwide. This data shows that
243 the CIET has sufficient reliability to be used as a research instrument for post-testing. No pre-testing is
244 being proposed in the current work. At the author's institution, classes in mechanical engineering run
245 double sections because classes are capped at 24 students. Hence the CIET will be administered to both
246 sections but only one section will be instructed using the proposed presentation. The primary purpose of
247 this paper to interest other thermodynamics instructors in the proposed presentation so that the assessment
248 can be performed with a large number of students, and the normal distribution can be used (not possible
249 with n=24) to assess the effect of the proposed approach on the population of students studying
250 thermodynamics .
251

252 **5. Conclusions:** A number of shortcomings in the classical presentation of the second law of
253 thermodynamics as found in contemporary engineering textbooks have been pointed out from a
254 pedagogical perspective. A new presentation that uses thought experiments about real rather than
255 imaginary processes to derive second-law results has been proposed. The proposed derivation has
256 conceptual implications. The effectiveness of the proposed presentation can therefore be measured using a
257 reliable concept inventory. The author urges fellow thermodynamics instructors to examine the problem
258 described, and consider an educational experiment with the proposed solution.

259

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